

Exam 1, part 2 - Solutions

Y (a) \leftarrow $F = mbv^{3/2}$ $\sum F = m\ddot{x}$ $-mbv^{3/2} = m\frac{dv}{dt}$ let $\frac{dv}{dt} = \frac{dv}{dx} \frac{dx}{dt} = \frac{dv}{dx} v$

$-bv^{3/2} = \frac{dv}{dx} v$

$\int_{x_0}^x -dx = \int_{v_0}^v \frac{v^{-1/2}}{b} dv$

$-(x-x_0) = \frac{2}{b} v^{1/2} \Big|_{v_0}^v$

$-(x-x_0) = \frac{2}{b} (\sqrt{v} - \sqrt{v_0})$

$v(x) = \left[\sqrt{v_0} - \frac{b}{2} (x-x_0) \right]^2$

(b) the farthest it gets is when $v=0$; i.e. it stops.

$0 = \sqrt{v_0} - \frac{b}{2} (x-x_0)$

so $(x-x_0) = \frac{2}{b} \sqrt{v_0}$

$$2/ \quad E = T + U$$

$$E = \frac{1}{2}mv^2 + \left(-\frac{x^2}{2} + \frac{x^4}{4} \right)$$

$$\frac{dE}{dt} = \frac{d}{dt} \left(\frac{1}{2}m\dot{x}^2 - \frac{x^2}{2} + \frac{x^4}{4} \right)$$

$$\frac{dE}{dt} = m\dot{x}\ddot{x} - x\dot{x} + x^3\dot{x} \quad \text{but note that } \sum F = m\ddot{x}$$

$$= m\dot{x} \left(\frac{x}{m} - \frac{x^3}{m} \right) - x\dot{x} + x^3\dot{x} \quad \begin{array}{l} +x - x^3 = m\ddot{x} \\ \frac{x}{m} - \frac{x^3}{m} = \ddot{x} \end{array}$$

$$= \dot{x}x - \dot{x}x^3 - x\dot{x} + x^3\dot{x}$$

$$\frac{dE}{dt} = 0 \quad \checkmark$$

Since energy does not change over time,

E is a constant + conserved.

3/ (a) determine the force... does it look like $F = -kx$? when x is a displacement from equilibrium.

$$\vec{F} = -\vec{\nabla}U$$

only in
x-direction

$$F_x = -\frac{\partial}{\partial x} \left[U_0 \left(1 - e^{-(x-x_0)/\delta} \right)^2 - U_0 \right]$$

$$= -U_0 (2) \left(1 - e^{-(x-x_0)/\delta} \right) (-1) \left(-\frac{1}{\delta} \right) e^{-(x-x_0)/\delta}$$

$$F_x = -\frac{2U_0}{\delta} \left(1 - e^{-(x-x_0)/\delta} \right) e^{-(x-x_0)/\delta}$$

for small separations $(x-x_0) \ll \delta$ use expansion of $e^x = 1 + x + \frac{x^2}{2!} + \dots$

$$F_x = -\frac{2U_0}{\delta} \left[1 - \left(1 - \frac{(x-x_0)}{\delta} + \frac{(x-x_0)^2}{2!\delta^2} + \dots \right) \right] \left(1 - \frac{(x-x_0)}{\delta} + \frac{(\quad)^2}{2!\delta^2} + \dots \right)$$

$$F_x = -\frac{2U_0}{\delta^2} (x-x_0) + \dots \text{ higher order } (x-x_0) \text{ terms}$$

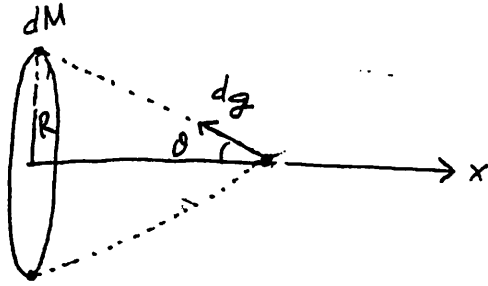
this is of the form $F = -kx$ where $k = \frac{2U_0}{\delta^2}$

(b) period of oscillation $T = \frac{2\pi}{\omega_0} = \frac{2\pi}{\sqrt{\frac{k}{m}}} = \frac{2\pi}{\sqrt{\frac{2U_0}{\delta^2 m}}}$

$$T = 2\pi \delta \sqrt{\frac{2m}{U_0}}$$

4/

(a)



note that only \hat{x} component will be non-zero once all contributions are added

$$g_x = - \int G \frac{dM}{r^2} \cos \theta$$

$$= - G \int \frac{dM}{(x^2 + R^2)} \frac{x}{(x^2 + R^2)^{1/2}}$$

note that for every dM , $x + R$ are the same!

$$= \frac{-Gx}{(x^2 + R^2)^{3/2}} \int dM$$

$$g_x = \frac{-GMx}{(x^2 + R^2)^{3/2}} \Rightarrow$$

$$\vec{g} = \frac{-GMx}{(x^2 + R^2)^{3/2}} \hat{x}$$

(b) can integrate $\int \vec{g} \cdot d\vec{r} = -\Delta U$ (since $\vec{g} = -\vec{\nabla} U$)

or do the potential directly!

$$\phi = - \int \frac{G dM}{r}$$

$$= - \int \frac{G dM}{(R^2 + x^2)^{1/2}}$$

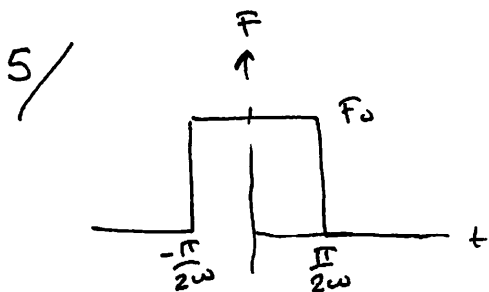
again, every dM has same R, x .

$$\phi = \frac{-GM}{(R^2 + x^2)^{1/2}}$$

$$4/(c) \quad \vec{F} = m \vec{g}$$

$$= m \left(-\frac{GMx}{(x^2 + R^2)^{3/2}} \right) \hat{x}$$

$$\vec{F} = -\frac{GMm x}{(x^2 + R^2)^{3/2}} \hat{x}$$



$$F(t) = \begin{cases} F_0 & -\frac{\pi}{2\omega} < t < \frac{\pi}{2\omega} \\ 0 & \text{other times in the period} \end{cases}$$

represent this as a Fourier series

$$F(t) = \frac{1}{2} a_0 + \sum_{n=1}^{\infty} [a_n \cos(n\omega t) + b_n \sin(n\omega t)]$$

$$\text{where } a_n = \frac{2}{\pi} \int F(t) \cos(n\omega t) dt$$

$$b_n = \frac{2}{\pi} \int F(t) \sin(n\omega t) dt$$

here, b_n 's = 0 since it is an even function. you can show this explicitly also ...

$$b_n = \frac{2}{\pi} \int_{-\pi/2\omega}^{+\pi/2\omega} F_0 \sin(n\omega t) dt$$

$$= \frac{2}{\pi} F_0 \left[-\frac{\cos(n\omega t)}{n\omega} \right]_{-\pi/2\omega}^{+\pi/2\omega}$$

$$b_n = \frac{-F_0}{n\pi} \left[\cos \frac{n\pi}{2} - \cos \left(-\frac{n\pi}{2} \right) \right] \quad \text{now, } \cos(\theta) = \cos(-\theta)$$

$$b_n = \frac{-F_0}{n\pi} \left[\cos \frac{n\pi}{2} - \cos \frac{n\pi}{2} \right] = 0$$

try the a's...

$$a_0 = \frac{1}{\pi} \int_{-\frac{\pi}{2\omega}}^{+\frac{\pi}{2\omega}} F_0 dt = \frac{1}{\pi} F_0 t \Big|_{-\frac{\pi}{2\omega}}^{+\frac{\pi}{2\omega}} = \frac{1}{\pi} F_0 \frac{\pi}{\omega} = F_0$$

$$a_n = \frac{1}{\pi} \int_{-\frac{\pi}{2\omega}}^{+\frac{\pi}{2\omega}} F_0 \cos(n\omega t) dt$$

$$= \frac{\omega F_0}{\pi n \omega} \sin(n\omega t) \Big|_{-\frac{\pi}{2\omega}}^{+\frac{\pi}{2\omega}}$$

$$= \frac{F_0}{\pi n} \left[\sin \frac{n\pi}{2} - \sin \left(-\frac{n\pi}{2} \right) \right] \quad \text{now } \sin(-\theta) = -\sin(\theta)$$

$$a_n = \frac{2F_0}{\pi n} \sin\left(\frac{n\pi}{2}\right)$$

zero when n is even
 $n = \pm 1$ when n is odd

$$\text{for } n=1 \quad a_1 = \frac{2F_0}{\pi}$$

$$n=3 \quad a_3 = -\frac{2F_0}{3\pi}$$

$$F(t) = \frac{1}{2} F_0 + \frac{2F_0}{\pi} \cos(\omega t) - \frac{2F_0}{3\pi} \cos(3\omega t) + \dots - \dots$$