

Exam 1 - Solutions

1/ a. $\gamma = \frac{1}{\sqrt{1 - (\frac{v}{c})^2}}$

$$= \frac{1}{\sqrt{1 - (\frac{\sqrt{3}c}{2c})^2}} = \frac{1}{(1 - \frac{3}{4})^{1/2}} = \frac{1}{(\frac{1}{4})^{1/2}} = (4)^{1/2}$$

$\gamma = 2$

b. here, $ct = (3 \cdot 10^8)(\frac{40}{3} \cdot 10^{-9}) = 4m.$

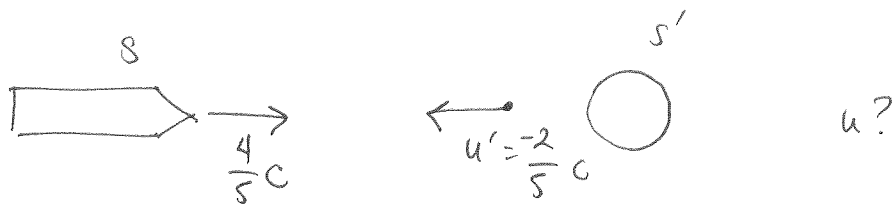
leave	$\begin{bmatrix} 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \end{bmatrix}$	
front expl	$\begin{bmatrix} 4m \\ 40m \end{bmatrix}$	$\begin{bmatrix} \text{see below} \\ \text{see below} \end{bmatrix}$	} use $\begin{bmatrix} ct' \\ x' \end{bmatrix} = \begin{bmatrix} \gamma & \gamma\beta \\ \gamma\beta & \gamma \end{bmatrix} \begin{bmatrix} ct \\ x \end{bmatrix}$
back expl	$\begin{bmatrix} 4m \\ 0 \end{bmatrix}$	$\begin{bmatrix} \text{see below} \\ \text{see below} \end{bmatrix}$	
	$\underbrace{\hspace{2cm}}_{S, Luis}$	$\underbrace{\hspace{2cm}}_{S', Josh}$	

front expl. $\begin{bmatrix} ct' \\ x' \end{bmatrix} = \begin{bmatrix} 2 & 2 \cdot \frac{\sqrt{3}}{2} \\ 2 \cdot \frac{\sqrt{3}}{2} & 2 \end{bmatrix} \begin{bmatrix} 4 \\ 40 \end{bmatrix} = \begin{bmatrix} 8 + 40\sqrt{3} \\ 4\sqrt{3} + 80 \end{bmatrix}$

back expl. $\begin{bmatrix} ct' \\ x' \end{bmatrix} = \begin{bmatrix} \cdot & \cdot \\ \cdot & \cdot \end{bmatrix} \begin{bmatrix} 4 \\ 0 \end{bmatrix} = \begin{bmatrix} 8 \\ 4\sqrt{3} \end{bmatrix}$

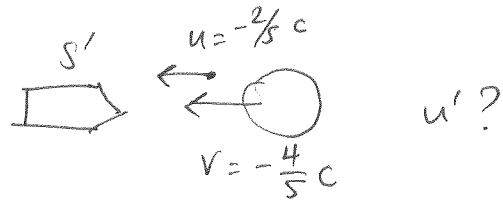
c. Ship leaves. Back explodes $\frac{8}{c}$ seconds later. An additional $\frac{40\sqrt{3}}{c}$ sec. later, the front explodes.

2/



transf. is $u' = \frac{u+V}{1 + \frac{uV}{c^2}}$

switch reference frames.



$$u' = \frac{-\frac{2}{5}c + \left(-\frac{4}{5}c\right)}{1 + \frac{\left(\frac{2}{5}c\right)\left(-\frac{4}{5}c\right)}{c^2}} = \frac{-\frac{4}{5}c}{1 + \frac{8}{25}} = \frac{-\frac{4}{5}c}{\frac{33}{25}}$$

$$= -\frac{30}{33}c = -\frac{10}{11}c$$

alternative soln. rearrange $u' = \frac{u+V}{1 + \frac{uV}{c^2}}$ to solve for u.

$$u = \frac{u' - V}{1 - \frac{u'V}{c^2}} = \frac{\left(-\frac{2}{5}c\right) - \left(\frac{4}{5}c\right)}{1 - \frac{\left(-\frac{2}{5}c\right)\left(\frac{4}{5}c\right)}{c^2}} = \text{same as above}$$

3/

energy conservation

$$E_i = E_f$$

momentum conserved

$$p_i = p_f$$

$$\gamma_{u_0} m_0 c^2 + \frac{hc}{\lambda} = \gamma_{u_f} m_f c^2$$

$$\gamma_{u_0} m_0 u_0 - \frac{h}{\lambda} = \gamma_{u_f} m_f u_f$$

4/ a. $E = \gamma_u mc^2$ and $p = \gamma_u mu$ when $\gamma_u = \frac{1}{\sqrt{1 - \frac{u^2}{c^2}}}$

to find u ...

$$\frac{p}{E} = \frac{\gamma_u mu}{\gamma_u mc^2} \rightarrow \frac{pc^2}{E} = u$$

subs. in values

$$u = \frac{(4E)c}{12E} = \left(\frac{c}{3} \right)$$

b. the invariant quantities are mass (m) and rest energy (mc^2)

5/ time order of events are fixed if the interval is timelike.

$$c\Delta t > \Delta x$$

here the 2 events are $\begin{bmatrix} 0 \\ 0 \end{bmatrix}, \begin{bmatrix} 240 \\ 200 \end{bmatrix}$ $\rightarrow ct = 3 \cdot 10^8 \cdot 800 \cdot 10^{-9} = 2400 \cdot 10^{-1} = 240$

so $240 > 200$

- It's time-like! So order of events are always the same.
- Note that the quantity $(c\Delta t)^2 - (\Delta x)^2$ is invariant. So no matter what reference frame is used, these two events will be time-like.
- The two events may be causally-related (but they don't have to be).

$$6/ (a) \quad E = hf = \frac{hc}{\lambda}$$

$$= \frac{(6.63 \cdot 10^{-34}) (3 \cdot 10^8)}{500 \cdot 10^{-9}} = 3.99 \cdot 10^{-19} \text{ J}$$

$$(b) \quad I = \frac{\text{power}}{\text{area}} = \frac{(hf) (dn/dt)}{\text{area}}$$

$$\frac{dn}{dt} = \frac{I \cdot \text{area}}{hf} \quad \text{where } f = \frac{c}{\lambda}$$

$$= \frac{(1.5 \cdot 10^3) (1 \cdot 10^{14})}{(6.63 \cdot 10^{-34}) \left(\frac{3 \cdot 10^8}{500 \cdot 10^{-9}} \right)} = 3.8 \cdot 10^{35} \frac{\text{photons}}{\text{sec}}$$

$$7/ \quad \Delta E = \Delta mc^2$$

$$\Delta m = \frac{\Delta E}{c^2}$$

$$\text{here } \Delta E = \text{power} \cdot \text{time} \rightarrow 1 \text{ s}$$

$$\rightarrow 1.5 \cdot 10^3 \frac{\text{J}}{\text{s area}} \cdot \text{area}$$

$$= \frac{(1.5 \cdot 10^3) (1) (1 \cdot 10^{14})}{(3 \cdot 10^8)^2}$$

$$= 1.7 \text{ kg in } 1 \text{ s.}$$

$$8/ \quad \phi = hf = \frac{hc}{\lambda}$$

$$\lambda = \frac{hc}{\phi} = \frac{(6.63 \cdot 10^{-34}) (3 \cdot 10^8)}{(1.2) (1.6 \cdot 10^{-19})} = 1.0 \cdot 10^{-6} \text{ m}$$

change eV \rightarrow J

$\Rightarrow \lambda$ must be equal to, or less than, $1 \cdot 10^{-6} \text{ m}$

9/ The Lorentz transformation is
$$\begin{bmatrix} ct' \\ x' \end{bmatrix} = \begin{bmatrix} \gamma & \gamma\beta \\ \gamma\beta & \gamma \end{bmatrix} \begin{bmatrix} ct \\ x \end{bmatrix}.$$

It will reduce the Galilean transformation in the limit of $\frac{v}{c} \ll 1$.

Since $\gamma = \frac{1}{\sqrt{1 - (\frac{v}{c})^2}} \rightarrow 1$ when $\frac{v}{c} \ll 1$.

$\beta = \frac{v}{c}$. Strictly speaking, $\beta \rightarrow 0$. However, let's keep its $\frac{v}{c}$ dependence, for now.

The transformation is then
$$\begin{bmatrix} ct' \\ x' \end{bmatrix} = \begin{bmatrix} 1 & \beta \\ \beta & 1 \end{bmatrix} \begin{bmatrix} ct \\ x \end{bmatrix}$$

$$ct' = ct + \beta x$$

$$x' = \beta ct + x.$$

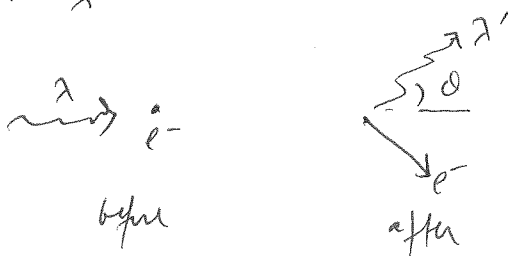
Sub in $\beta = \frac{v}{c}$, and simplify a little

$$\begin{bmatrix} t' = t + \frac{v}{c^2}x \\ x' = vt + x \end{bmatrix}$$

$t' = t$ in limit that $\frac{v}{c} \ll 1$,
However, if x is v. large, this isn't the case.
→ clearly Galilean.

10/ the Compton experiment is the scattering of light off free e^- , as shown below. His results show that the scattered light has a wavelength that depends on the scattering angle, $\lambda' - \lambda = \frac{h}{m_e c} (1 - \cos \theta)$.

The results may be understood w/ the photon framework, when light has momentum $p = \frac{h}{\lambda}$ & interacts w/ the e^- like a particle.



11/ Black body radiation is the emission of light from blackbodies. The spectrum of this emission has a peak, but would not be explained classically. In fact, classically, the spectrum should diverge at low wavelengths. (see picture). The problem was resolved by postulating that the energy of light is quantized, $E = nhf$, as opposed to being continuous.

