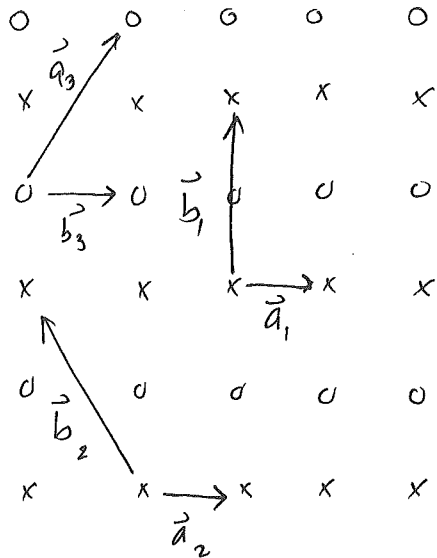
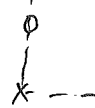


# Exam 1 - Solutions

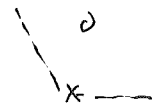
1/ (a) two examples of primitive cells



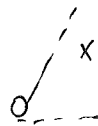
$\vec{a}_1, \vec{b}_1$  has the basis



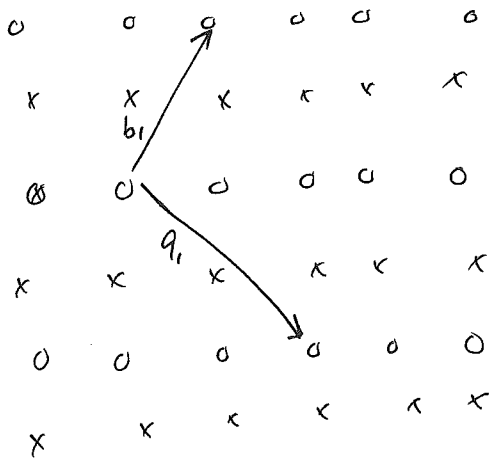
$\vec{a}_2, \vec{b}_2$  has the basis



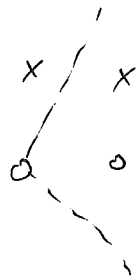
$\vec{a}_3, \vec{b}_3$  has the basis



(b) example of non-primitive cells

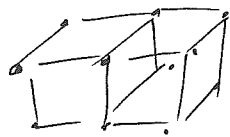


$\vec{a}_1, \vec{b}_1$  has basis



2) (a) Negative (-). Negative energy indicates that energy has to be added to unbind (or free) the ion from the crystal  
 $E \geq 0$

(b) 
$$-\frac{6e^2}{4\pi\epsilon_0 a^2}$$



6: the 6 nearest neighbors

a: that are a distance a away

$-e^2$ : the neighbor has the opposite charge

ex:  $\text{Na}^+$  surrounded by 6  $\text{Cl}^-$  ions

$$\frac{+12e^2}{4\pi\epsilon_0 a\sqrt{2}}$$

12: the 12 atoms

$a\sqrt{2}$ : that are a distance  $a\sqrt{2}$  away

$+e^2$ : they are the same charge

ex: the 8 closest  $\text{Na}^+$  ions to a single  $\text{Na}^+$

$$\frac{-8e^2}{\dots a\sqrt{3}}$$

8: the 8 atoms

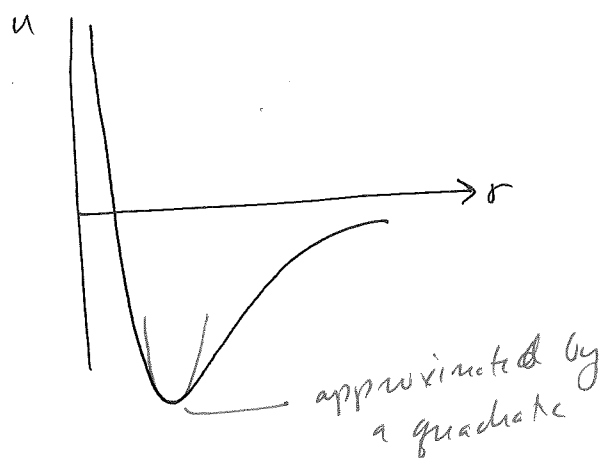
$a\sqrt{3}$ : that are a distance  $a\sqrt{3}$  away

$(-e^2)$ : they are opp. charge

Ex: the second closest  $\text{Cl}^-$  atoms to  $\text{Na}^+$

3/ (a)  $+\left(\frac{c}{r}\right)^2$  is the repulsion  
 $-\left(\frac{b}{r}\right)$  attraction.

(b) For a system to exhibit SHO, its  $U(x) = \frac{1}{2} kx^2$ .  
 In general, any system with a minimum in potential will show SHO in the limit of small displacements from equilibrium. That is



Show this using a Taylor series expansion around  $r = r_m$

$$U(r) \approx \underbrace{U(r_m)}_{\text{offset to potential}} + \underbrace{U'(r_m)(r-r_m)}_{=0 \text{ since } r_m \text{ is the minimum}} + \underbrace{\frac{U''(r_m)}{2!}(r-r_m)^2 + \dots}_{\text{the SH term.}}$$

$$U(r) = a \left[ c^2 r^{-2} - b r^{-1} \right]$$

$$U'(r) = a \left[ c^2 (-2) r^{-3} - b (-1) r^{-2} \right]$$

$$U''(r) = a \left[ c^2 (-2)(-3) r^{-4} - b (-1)(-2) r^{-3} \right]$$

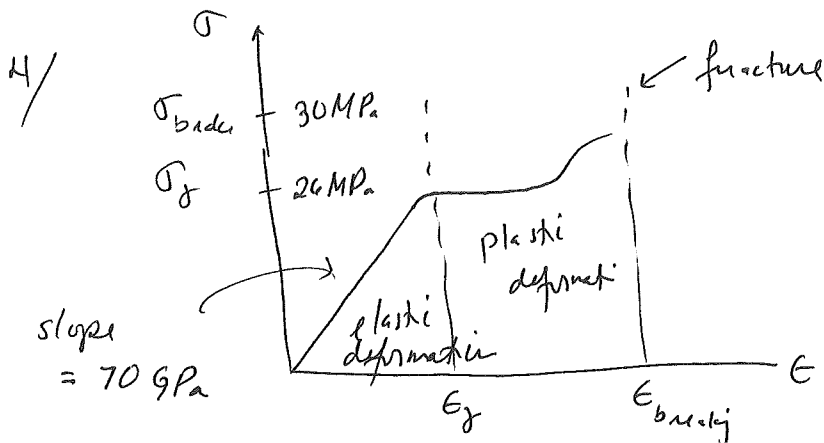
$$= a \left[ 6c^2 r^{-4} - 2b r^{-3} \right]$$

$$U''(r_m) = a^2 \left[ \frac{3c^2}{(2c^2/b)^4} - \frac{b}{(2c^2/b)^3} \right]$$

3 (cont'd)

$$u''(r_m) = \frac{ab^4}{8c^6}$$

$$\text{so } u(r) = u(r_m) + \frac{1}{2} \left[ \frac{ab^4}{8c^6} \right] [r - r_m]^2$$



$$Y = \frac{\sigma}{\epsilon}$$

$$\epsilon = \frac{\sigma}{Y} = \frac{F/A}{Y} = \frac{8500}{\pi(r_o^2 - r_i^2) \cdot 70 \cdot 10^9}$$

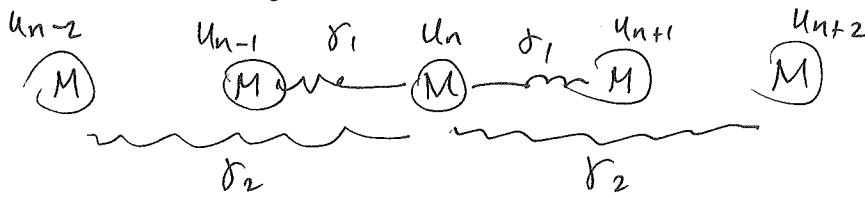
$$\epsilon = \frac{8500}{\pi \left[ \left( \frac{0.1425}{2} \right)^2 - \left( \frac{0.16}{2} \right)^2 \right] \cdot 70 \cdot 10^9} = 0.19 \cdot 10^{-3}$$

actual compresni

$$\frac{\Delta l}{l} = \epsilon$$

$$\Delta l = \epsilon l = 0.6 \text{ mm}$$

5/ Start by applying Newton's 2<sup>nd</sup> law to the  $n^{\text{th}}$  atom



$$F = m \frac{d^2 u_n}{dt^2}$$

$$-\delta_1 (u_n - u_{n-1}) + \delta_1 (u_{n+1} - u_n) - \delta_2 (u_n - u_{n-2}) + \delta_2 (u_{n+2} - u_n) = m \frac{d^2 u_n}{dt^2} \quad (1)$$

Propose the following solution

$$u_n(t) = u_0 e^{i(kan - \omega t)} \quad \text{for the } n^{\text{th}} \text{ atom} \quad (2a)$$

$$u_{n-1}(t) = u_0 e^{i(ka(n-1) - \omega t)} \quad (n-1)^{\text{th}} \text{ atom} \quad (2b)$$

and so on.

Sub in (2a), (2b) and similar eqns for  $n-2, n+2, n+1$  into (1).

In order for (2) to be a solution of (1),

$\omega$  will have to be related to  $k$  in some way.

This relationship,  $\omega(k)$ , is the dispersion relation.

6/ The group velocity  $V_g = \frac{d\omega}{dk}$ .

(a) In the limit  $k \rightarrow 0$ , we can approximate  $\cos ka \approx 1 - \frac{(ka)^2}{2}$

$$\omega^2 \approx \frac{2\gamma}{M} \left[ 1 - \left( 1 - \frac{(ka)^2}{2} \right) \right]$$

$$\omega^2 \approx \frac{2\gamma}{M} \left[ \frac{(ka)^2}{4} \right]$$

$$\omega = \sqrt{\frac{\gamma}{2M}} ka$$

$$\frac{d\omega}{dk} = \sqrt{\frac{\gamma}{2M}} a$$

(b) In the limit  $k \rightarrow \frac{\pi}{a}$ , take the derivative directly

$$\omega = \sqrt{\frac{2\gamma}{M}} [1 - \cos(ka)]^{1/2}$$

$$\frac{d\omega}{dk} = \sqrt{\frac{2\gamma}{M}} \left( \frac{1}{2} \right) (1 - \cos(ka))^{-1/2} (+ \sin(ka)) (a)$$

$$\frac{d\omega}{dk} = \sqrt{\frac{\gamma}{2M}} \frac{a \sin(ka)}{[1 - \cos(ka)]^{1/2}}$$

$$\left. \frac{d\omega}{dk} \right|_{k=\frac{\pi}{a}} = \sqrt{\frac{a \sin(\pi)}{[1 - \cos(\pi)]^{1/2}}} = 0$$

\* note: you can also use this to evaluate  $\left. \frac{d\omega}{dk} \right|_{k=0}$ , but

you'll note that this leads to  $\frac{0}{0}$  if you just evaluate at  $k=0$ .

See your calc text on how to do this.

6 cont'd/

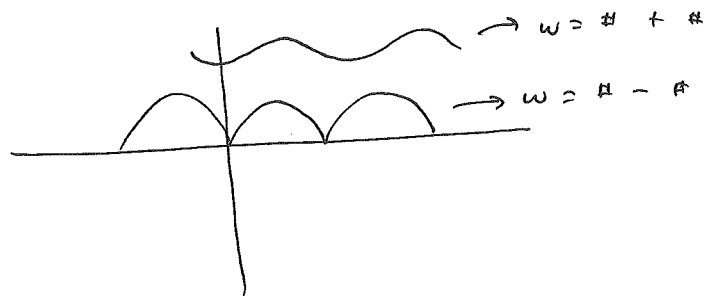
(c) since  $k = \frac{2\pi}{\lambda}$ ,  $k \rightarrow 0$  is the long wavelength limit

(d) there is no optical branch.

If there are 2+ branches, it means that  $\omega$  has 2+ possible positive values for a given  $k$ .

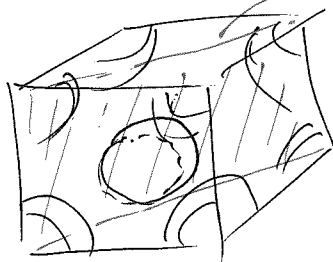
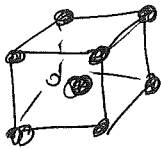
•  $\omega = \pm \pm$

↑ indicates 2 values of  $\omega$

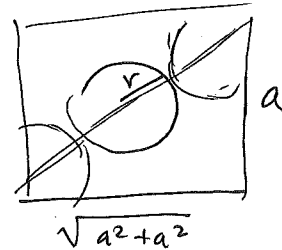


7/

bcc



along the central diagonal



packing fraction =  $\frac{\text{Volume ions}}{\text{Volume cube}}$

$$= \frac{(1 \times \frac{4}{3} \pi r^3) + 8 \times \frac{1}{8} (\frac{4}{3} \pi r^3)}{a^3}$$

$$a^2 + a^2 + a^2 = (4r)^2$$

$$3a^2 = 16r^2$$

$$\frac{\sqrt{3}}{4} a = r$$

$$= \frac{\frac{8}{3} \pi r^3}{a^3} = \frac{\frac{8}{3} \pi \left[ \frac{\sqrt{3}}{4} a \right]^3}{a^3} = \frac{\frac{8}{3} \pi \sqrt{3}}{4 \cdot 4 \cdot 4} = \frac{\sqrt{3} \pi}{8} = 68\%$$

- 8 / (a) Li, Mg, P - elements in same column have same # valence e<sup>-</sup>  
 increasing # in valence →  
 (b) Be, O, Ne - in same row, ionization energy increases w/ column →  
 (c) 1s<sup>6</sup> - all e<sup>-</sup> would be in the first state

9 / (a)  $2d \sin \theta = m\lambda$

$$d = \frac{m\lambda}{2 \sin \theta} = \frac{(1) (0.109 \cdot 10^{-9})}{2 (\sin 12.6^\circ)} = 0.25 \text{ nm}$$

(b)  $E = \frac{p^2}{2m}$  and  $\lambda = \frac{h}{p}$

$$\lambda = \frac{h}{\sqrt{2mE}} = \frac{6.6 \cdot 10^{-34}}{\sqrt{2 (9.1 \cdot 10^{-31}) (200 \cdot 10^3 \text{ eV} \cdot \frac{1.6 \cdot 10^{-19} \text{ J}}{\text{eV}})}}$$

$$= 2.7 \cdot 10^{-12} \text{ m}$$

$$= 0.0027 \text{ nm}$$